

Remote Synchronization Reveals Network Symmetries and Functional Modules

Vincenzo Nicosia,¹ Miguel Valencia,² Mario Chavez,³ Albert Díaz-Guilera,⁴ and Vito Latora^{1,5}

¹*School of Mathematical Sciences, Queen Mary University of London, London E1 4NS, United Kingdom*

²*Neurophysiology Laboratory, CIMA, University of Navarra, 31008 Pamplona, Spain*

³*CNRS UMR-7225, Hôpital de la Salpêtrière, 47 Boulevard de l'Hôpital, 75013 Paris, France*

⁴*Department de Física Fonamental, Facultat de Física, Universitat de Barcelona, E-08028 Spain*

⁵*Dipartimento di Fisica ed Astronomia, Università di Catania and INFN, I-95123 Catania, Italy*

(Received 26 November 2012; revised manuscript received 12 February 2013; published 25 April 2013)

We study a Kuramoto model in which the oscillators are associated with the nodes of a complex network and the interactions include a phase frustration, thus preventing full synchronization. The system organizes into a regime of remote synchronization where pairs of nodes with the same network symmetry are fully synchronized, despite their distance on the graph. We provide analytical arguments to explain this result, and we show how the frustration parameter affects the distribution of phases. An application to brain networks suggests that anatomical symmetry plays a role in neural synchronization by determining correlated functional modules across distant locations.

DOI: [10.1103/PhysRevLett.110.174102](https://doi.org/10.1103/PhysRevLett.110.174102)

PACS numbers: 05.45.Xt, 87.19.lm, 89.75.Fb, 89.75.Kd

Synchronization of coupled dynamical units is a ubiquitous phenomenon in nature [1]. Remarkable examples include phase locking in laser arrays, rhythms of flashing fireflies, wave propagation in the heart, and also normal and abnormal correlations in the activity of different regions of the human brain [2–5]. In 1975 Y. Kuramoto proposed a simple microscopic model to study collective behaviors in large populations of interacting elements [6]. In its original formulation the Kuramoto model describes each unit of the system as an oscillator which continuously readjusts its frequency in order to minimize the difference between its phase and the phase of all the other oscillators. This model has shown very successful in understanding the spontaneous emergence of synchronization and, over the years, many variations have been considered [7–9]. Recently, the Kuramoto model has been also extended to sets of oscillators coupled through complex networks [2,10,11], and it has been found that the topology of the interaction network has a fundamental role in the emergence and stability of synchronized states [12,13]. In particular, the presence of communities—groups of tightly connected nodes—has a relevant impact on the path to synchronization [14–18], and units that are close to each other on the network, or belong to the same module or community [19], have a higher chance to exhibit similar dynamics. This implies that nodes in the same structural module share similar functions, which is a belief often supported by empirical findings [3,20]. However, various examples are found in nature where functional similarity is instead associated with morphological symmetry. In these cases, units with similar roles, which could potentially swap their position without altering the overall functioning of the system, appear in remote locations of the network. Some examples include cortical areas in brains [21], symmetric

organs in plants and vertebrates [22,23], and even atoms in complex molecules [24]. Therefore, identifying the sets of symmetric units of a complex system might be helpful to understand its organization. Finding the global symmetries in a graph, i.e., constructing its automorphism group, is a classical problem in graph theory. However, it is still unknown if this problem is polynomial or NP-complete [25,26], even if there exist polynomial-time algorithms for graphs with bounded maximum degree [27]. Recent works have focused instead on defining and detecting local symmetries in complex networks [28,29]. Nevertheless, the interplay between the structural symmetries of a network and the dynamics of processes occurring over the network has been studied only marginally [30–32], or for specific small network motifs [33–35].

In this Letter we show that network symmetries play a central role in the synchronization of a system. We consider networks of identical Kuramoto oscillators, in which a phase frustration parameter forces connected nodes to maintain a finite phase difference, thus hindering the attainment of full synchronization. We prove that the configuration of phases at the synchronized state reflects the symmetries of the underlying coupling network. In particular, two nodes with the same symmetry have identical phases, i.e., are fully synchronized, despite the distance between the two nodes on the graph. Such remote synchronization behavior is here induced by the network symmetries and not by an initial *ad hoc* choice of different natural frequencies [30].

Let us consider N identical oscillators associated to the nodes of a connected graph $G(\mathcal{N}, \mathcal{L})$, with $N = |\mathcal{N}|$ nodes and $K = |\mathcal{L}|$ links. Each node i is characterized, at time t , by a phase $\theta_i(t)$ whose time evolution is governed by the equation

$$\dot{\theta}_i = \omega + \lambda \sum_{j=1}^N a_{ij} \sin(\theta_j - \theta_i - \alpha). \quad (1)$$

Here ω is the natural frequency, identical for all the oscillators, and $A \equiv \{a_{ij}\}$ is the adjacency matrix of the coupling graph. The model has two control parameters: $\lambda > 0$ accounting for the strength of the interaction, and α , the phase frustration parameter ranging in $[0, \pi/2]$. When $\alpha = 0$, the model reduces to a network of identical Kuramoto oscillators. In this case, the fully synchronized state is globally stable for a set of initial conditions having finite measure [6,9] and the transient dynamics closely reflects the structure of the graph, so that nodes belonging to the same structural module evolve similarly in time [14]. However, the synchronized state can coexist with other nontrivial attractors, e.g., uniformly twisted waves, especially if the coupling topology is regular and sparse (see Ref. [12] for a discussion about the size of the sync basin). Instead, if the oscillators are not identical the frequency distribution tends to separate their phases and, as a result, there is a transition from an incoherent state (with order parameter $r = \frac{1}{N} |\sum_{j=1}^N e^{i\theta_j}|$ equal to 0) to a synchronized one ($r \neq 0$) at a critical value λ_c of the coupling strength.

The introduction of a phase frustration $\alpha \neq 0$ forces directly connected oscillators to maintain a constant phase difference [36]. In particular, we found that for a wide range of $\alpha > 0$ the dynamics in Eq. (1) reaches a stationary state in which the oscillators at two symmetric nodes have exactly the same phase, and this phase differs from the phases of nodes with different symmetries. Let us first illustrate this behavior and the effect of α on the three graphs G_a , G_b and G_c shown in Fig. 1. In the three topmost panels of Fig. 2 we report the results of the numerical integration of Eq. (1) on the graph G_a for three different values of α . We find that, after a transient, the system settles into a stationary state in which, at any time t , the phases are grouped into four different trajectories: $\theta_1(t)$, $\theta_2(t) = \theta_3(t)$, $\theta_4(t) = \theta_7(t)$ and $\theta_5(t) = \theta_6(t)$. In general, by increasing α up to a certain value α_c we better separate the four trajectories.

The four clusters of nodes obtained for $\alpha < \alpha_c$ are identified by a color code in Fig. 1. We notice that each cluster groups together all the nodes with the same symmetry. In this way two distant nodes of the graph, e.g., node 4 and node 7, are fully synchronized even if the other nodes in the paths connecting them have different phases. In this respect, what we observe is a remote synchronization [30]. We have found similar results for the linear chain and for the Bethe lattice (see nodes with the same colors in G_b and G_c in Fig. 1).

Notice that if the system reaches a synchronized state and α is small enough, Eq. (1) can be linearized by replacing the sinus with its argument. We obtain

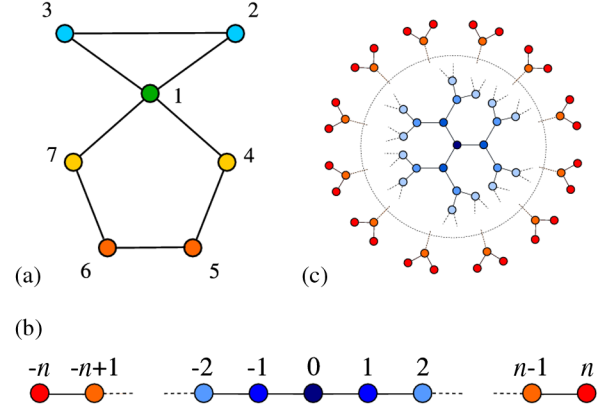


FIG. 1 (color online). The presence of frustration reveals clusters of symmetric nodes. The color code represents the phases of nodes at a given time in the stationary state. (a) In the first graph (G_a), node 2 is synchronized to node 3, node 4 to node 7, and node 5 to node 6. (b) In a finite chain (G_b), pairs of nodes symmetrically placed with respect to the central node are perfectly synchronized. (c) In a finite Bethe lattice (G_c) all the nodes placed at the same distance from the center have equal phases.

$$\dot{\theta}_i = \omega - \lambda \left[\sum_{j=1}^N L_{ij} \theta_j + \alpha k_i \right] \quad (2)$$

where $k_i = \sum_j a_{ij}$ is the degree of node i , L_{ij} are the entries of the Laplacian matrix of the graph $L \equiv D - A$, and D is a diagonal matrix such that $D_{ii} = k_i$. Without loss of generality, we can set $\lambda = 1$, $\omega = 0$. If the system is synchronized then $\theta_i = \Omega$, $\forall i$, so that the phases must satisfy the equations $\sum_{j=1}^N L_{ij} \theta_j = \alpha[\langle k \rangle - k_i]$ at any time, or equivalently,

$$L\theta = \alpha[\langle k \rangle \mathbf{1} - \mathbf{k}] \quad (3)$$

where $\langle k \rangle = N^{-1} \sum_i k_i$ is the average degree of the network. This corresponds to a synchronization frequency $\theta_i = \Omega = -\alpha \langle k \rangle \forall i$. In a connected graph the Laplacian matrix has one null eigenvalue and the system of Eq. (3) is singular. Consequently, at each time t we can solve the system by computing the phase difference between each node and a given node chosen as reference. For instance, if in G_a we define $\phi_j(t) = \theta_j(t) - \theta_1(t)$, $j = 2, \dots, 7$, by solving Eq. (3) we obtain $\phi_2 = \phi_3 = \alpha[\langle k \rangle - 2]$, $\phi_4 = \phi_7 = 2\alpha[\langle k \rangle - 2]$ and $\phi_5 = \phi_6 = 3\alpha[\langle k \rangle - 2]$. This is in agreement with the results of the simulations: the phases are clustered into four groups, with nodes with the same symmetry having the same phase, and nodes with different symmetries being separated by a phase lag that depends on α as in the relations found above. An analogous analytical expression can be derived for a finite chain (graph G_b in Fig. 1), for which we obtain $\theta_n - \theta_{n-i} = \theta_{-n} - \theta_{-n+i} = [(i(i+1)/2)\langle k \rangle - i^2]\alpha$ and $\theta_n - \theta_0 = \theta_{-n} - \theta_0 = [(n(n+1)/2)\langle k \rangle - n^2]\alpha$. Consequently, two

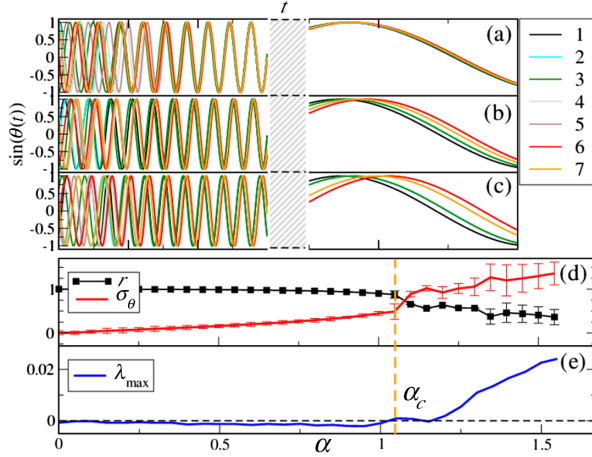


FIG. 2 (color online). The figure refers to the coupling topology G_a in Fig. 1(a). Panels (a)–(c): after an initial transient the system reaches a phase-locked synchronized state in which symmetric nodes have the same phase. The panels correspond to three different values of the frustration parameter, respectively, (a) $\alpha = 0.1$, (b) $\alpha = 0.5$, (c) $\alpha = 0.8$. Panel (d): for $\alpha > \alpha_c$ the synchronized state becomes unstable, the order parameter decreases while the dispersion of the phases σ_θ increases. Panel (e): the maximum Lyapunov exponent λ_{\max} of oscillators coupled through G_a becomes positive for $\alpha > \alpha_c$, and the systems enters a chaotic regime. The dashed yellow line indicates the approximate position of α_c .

nodes symmetrically placed with respect to node 0 will have identical phases.

We now provide a general argument to explain why the synchronization of Eq. (1) is related to graph symmetries. A graph $G(\mathcal{N}, \mathcal{L})$ has a symmetry if and only if it is possible to find a bijection $\pi: \mathcal{N} \rightarrow \mathcal{N}$ which preserves the adjacency relation of G , i.e., which is an automorphism for G . Formally, this means that there exists a permutation matrix $P = P(\pi)$ such that $PAP^{-1} = A$. If P corresponds to an automorphism of G then P commutes with A , i.e. $PA = AP$, and PAP^{-1} performs a relabeling of the nodes of the original graph which preserves the adjacency matrix [37]. In general a graph can admit more than one automorphism. For instance, graph G_a in Fig. 1 has at least three nontrivial bijections which preserve the adjacency matrix, namely,

$$\begin{aligned}\pi_1: (1, 2, 3, 4, 5, 6, 7) &\rightarrow (1, 3, 2, 4, 5, 6, 7) \\ \pi_2: (1, 2, 3, 4, 5, 6, 7) &\rightarrow (1, 2, 3, 7, 6, 5, 4) \\ \pi_3: (1, 2, 3, 4, 5, 6, 7) &\rightarrow (1, 3, 2, 7, 6, 5, 4).\end{aligned}$$

Node 2 and node 3 are symmetric because we can relabel the nodes of G_a (e.g., by means of either π_1 or π_3) so that node 2 is mapped into node 3 and vice versa, and the adjacency matrix of G_a is left unchanged. Similarly, for the pairs $\{4, 7\}$ and $\{5, 6\}$, there are two different relabelings which preserve adjacency relations, i.e., π_2 and π_3 . In terms of symmetries, the graph G has the following four

different classes of nodes: $C_1 = \{1\}$, $C_2 = \{2, 3\}$, $C_3 = \{4, 7\}$, $C_4 = \{5, 6\}$. Now, if a permutation of the nodes is an automorphism of G , then $PLP^{-1} = PDP^{-1} - PAP^{-1} = D - A = L$; i.e., the associated permutation matrix P also commutes with the Laplacian matrix of the graph. By left-multiplying both sides of Eq. (3) by P , we get $PL\theta = \alpha P[\langle k \rangle \mathbf{1} - \mathbf{k}]$. Since $PL = LP$ (P commutes with L) and $P\mathbf{k} = \mathbf{k}$ (symmetric nodes have the same degree) then we have

$$LP\theta = \alpha[\langle k \rangle \mathbf{1} - \mathbf{k}] \quad (4)$$

Combining Eqs. (3) and (4), we finally obtain the linear system,

$$LP\theta = L\theta, \quad (5)$$

which is singular, i.e., has one free variable. Again, it can be solved by leaving free one of the N variables θ_i , setting $\phi_j = \theta_j - \theta_i$ and considering the new system $\tilde{L}\tilde{P}\phi = \tilde{L}\phi$. The matrix \tilde{P} is obtained from P by removing the row and the column corresponding to node i . If P does not permute node i with another node, then \tilde{P} is still a permutation matrix. Similarly, \tilde{L} is the reduced Laplacian, i.e. the matrix obtained from the Laplacian by deleting the i -th row and the i -th column. By left-multiplying by \tilde{L}^{-1} , which is not singular, we obtain

$$\tilde{P}\phi = \phi. \quad (6)$$

Since $\tilde{P}\phi$ is a permutation of the phases of symmetric nodes, Eq. (6) implies that the phases of symmetric nodes will be equal at any time, whereas by solving Eq. (4) we can get the values of the corresponding phases. This argument is valid for small values of α , since the linearization of Eq. (1) is possible only if $\sin(x - \alpha) \simeq (x - \alpha)$, but as shown in Figs. 2(a)–2(c), we observe the formation of the same perfectly synchronized clusters of symmetric nodes for a wide range of α . However, when α becomes larger than a certain value α_c , the assumption $\dot{\theta}_i = \Omega$, $\forall i$ does not hold any more and the global synchronized state loses stability. By looking at Figs. 2(d) and 2(e) we notice that for $\alpha > \alpha_c$, with $\alpha_c \simeq 1.05$ for the graph G_a , the value of r steadily decreases while the dispersion of phases increases, until it reaches the expected value $\bar{\sigma}_\theta \simeq 1.39$ for a system of seven incoherent oscillators (see Fig. 2(d) and Ref. [38]). Moreover, for $\alpha > \alpha_c$ the maximal Lyapunov exponent of the system λ_{\max} becomes positive and the system enters a chaotic regime [see Fig. 2(e)]. Interestingly, the results reported in Figs. 3(a)–3(d) confirm that in this regime the coherence of symmetric nodes, measured by the pairwise order parameter r_2 , is higher than expected for incoherent oscillators (refer to Ref. [38] for additional details). Figure 3(e) shows that for $\alpha > \alpha_c$ the system exhibits metastable, partially synchronized states, in which pairs of symmetric nodes alternates intervals of perfect synchronization with intervals of complete incoherence. We point out that in this regime chimera states

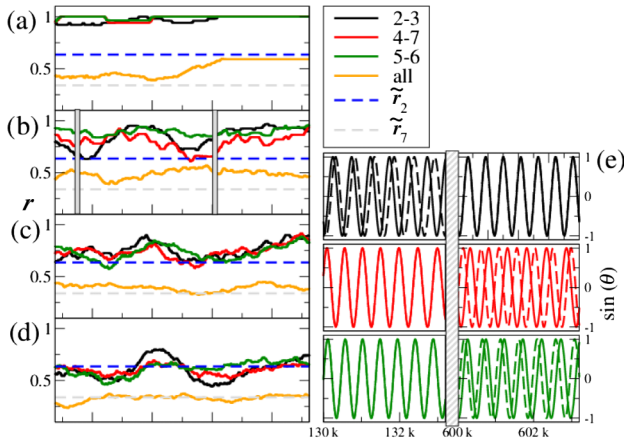


FIG. 3 (color online). Chaotic regime in G_a when $\alpha > \alpha_c$. Panel (a)–(d): running averages of r (orange) and pairwise order parameters r_2 (black, red and green lines) for typical trajectories of oscillators coupled through graph G_a , when α is, respectively, equal to (a) 1.3, (b) 1.4, (c) 1.5 and (d) 1.55. The dashed lines indicate the expected synchronization level for a system of two (blue line, $\bar{r}_2 = 2/\pi$) and seven incoherent oscillators (gray line, $\bar{r}_7 \approx 0.338\dots$). Panel (e): the plot of the phases of pairs of symmetric nodes for $\alpha = 1.4$ in two different temporal intervals (the shaded gray regions in panel (b)) reveal the existence of metastable, partially synchronized states.

could potentially occur [39–42] and could even coexist with remote synchronization for $\alpha < \alpha_c$. Qualitatively similar results are obtained for different coupling topologies, but the actual value of α_c seems to depend on the structure of the coupling network in a non-trivial way.

Application to the brain.—As an example, we investigate here the role of symmetry in the human brain by considering anatomical and functional brain connectivity graphs defined on the same set of $N = 90$ cortical areas (see details in Ref. [38]). We have first constructed a graph of anatomical brain connectivity as obtained from DW-MRI data [43], where links represent axonal fibers, and we used this graph as a backbone network to integrate Eq. (1). We identified candidate pairs of anatomically symmetric areas by means of agglomerative clustering, i.e. grouping together nodes having close phases at the stationary state (full dendrogram and details are provided in Ref. [38]). Then, we considered the graph of functional brain connectivity, in which links represent statistically significant correlations between the BOLD fMRI time-series of cortical areas (see details in Ref. [38]). Figure 4 illustrates the results for $\alpha = 0.5$ (we obtained qualitatively similar results in a wide range of α). Consider nodes 57 and 74, corresponding respectively to the green and blue areas in panel (a). Not only the two areas are spatially separated, but there is no edge connecting the two corresponding nodes in the anatomical connectivity network. However, the two nodes are detected as a candidate symmetric pair since at the stationary state of the Kuramoto

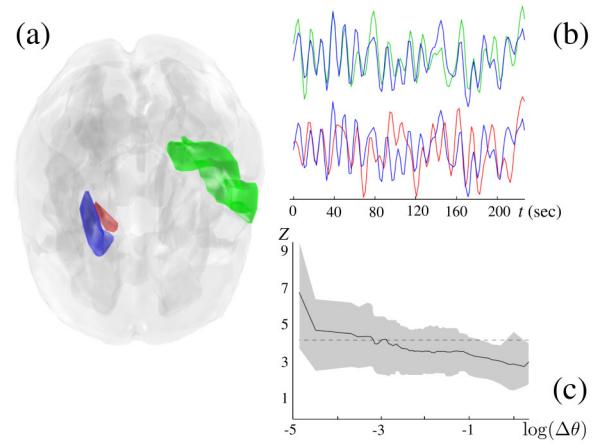


FIG. 4 (color online). (a) Brain areas with similar and dissimilar phases of the frustrated Kuramoto model are colored and superimposed onto an anatomical image. (b) Examples of functional data from one subject recorded at the brain areas indicated in panel (a). Colors are the same as those used in the anatomical image. (c) Functional correlation Z between pairs of nodes as a function of their phase differences $\Delta\theta$ according to the simulated Kuramoto dynamics. The black solid curve corresponds to the average value over all the subjects, while the gray area covers the 5th and the 95th percentiles of the distribution. The dashed horizontal line indicates the threshold for statistical significant correlations ($p < 0.05$, corrected for multiple comparisons).

dynamics in Eq. (1) the oscillators associated to these two nodes have very close phases (see dendrogram in Ref. [38]). As shown in Fig. 4(b), the BOLD fMRI signals corresponding to nodes 57 and 74 also are strongly synchronized. We obtain remarkably different results when we consider node 74 and node 76. These nodes correspond to two spatially adjacent areas of the brain [the red and blue regions in Fig. 4(a)] and are directly connected in the anatomical connectivity network. However, at the stationary state of Eq. (1) the phase difference of the oscillators associated to node 74 and 76 is quite large. Interestingly, in this case the fMRI time-series associated to these nodes are much less similar to each other [see the two bottom trajectories reported in Fig. 4(b)].

To quantify this effect, we plot in Fig. 4(c) the average functional correlation Z between the fMRI activity of pairs of brain areas as a function of the phase differences $\Delta\theta$ between the phases of the corresponding oscillators, obtained from the dynamics of Eq. (1) on the anatomical connectivity network. The fact that Z decreases with $\Delta\theta$ suggests that structural symmetry plays an important role in determining human brain functions. In fact, the functional activities of anatomically symmetric areas can be strongly correlated, even if the areas are distant in space. These results suggest that the study of anatomical symmetries in neural systems might provide meaningful insights about the functional organization of distant neural assemblies during diverse cognitive or pathological states [21]. Applied to other connectivity networks as a method to spot

potential network symmetries, our study could provide new insights on the interplay between structure and dynamics in complex systems.

The authors thank Yasser Iturria-Medina for sharing the DTI connectivity data used in the study and Simone Severini for useful comments. M. V. acknowledges financial support from the Spanish Ministry of Science and Innovation, Juan de la Cierva Programme Ref. No. JCI-2010-07876. A.D.-G. acknowledges support from the Spanish DGICYT Grant No. FIS2009-13730 and from the Generalitat de Catalunya No. 2009SGR00838. This work was supported by the EU-LASAGNE Project, Contract No. 318132 (STREP).

-
- [1] A. Pikovsky, M. Rosenblum, and J. Kurths, *Synchronization: a Universal Concept in Nonlinear Sciences* (Cambridge University Press, Cambridge, England, 2003).
 - [2] S. H. Strogatz, *Nature (London)* **410**, 268 (2001).
 - [3] E. Bullmore and O. Sporns, *Nat. Rev. Neurosci.* **10**, 186 (2009).
 - [4] M. Chavez, M. Valencia, V. Navarro, V. Latora, and J. Martinerie, *Phys. Rev. Lett.* **104**, 118701 (2010).
 - [5] P. Tass, *Phase Resetting in Medicine and Biology* (Springer, New York, 1999).
 - [6] Y. Kuramoto, *Lect. Notes Phys.* **39**, 420 (1975).
 - [7] H. Sakaguchi and Y. Kuramoto, *Prog. Theor. Phys.* **76**, 576 (1986).
 - [8] S. H. Strogatz, *Physica (Amsterdam)* **143**, 1 (2000).
 - [9] J. A. Acebrón, L. L. Bonilla, C. J. Pérez-Vicente, F. Ritort, and R. Spigler, *Rev. Mod. Phys.* **77**, 137 (2005).
 - [10] D. J. Watts and S. H. Strogatz, *Nature (London)* **393**, 440 (1998).
 - [11] S. Boccaletti, V. Latora, Y. Moreno, M. Chavez, and D.-U. Hwang, *Phys. Rep.* **424**, 175 (2006).
 - [12] D. A. Wiley, S. H. Strogatz, and M. Girvan, *Chaos* **16**, 015103 (2006).
 - [13] A. Arenas, A. Díaz-Guilera, J. Kurths, Y. Moreno, and C. Zhou, *Phys. Rep.* **469**, 93 (2008).
 - [14] A. Arenas, A. Díaz-Guilera, and C. J. Pérez-Vicente, *Phys. Rev. Lett.* **96**, 114102 (2006).
 - [15] S. Boccaletti, M. Ivanchenko, V. Latora, A. Pluchino, and A. Rapisarda, *Phys. Rev. E* **75**, 045102(R) (2007).
 - [16] J. Gómez-Gardeñes, Y. Moreno, and A. Arenas, *Phys. Rev. Lett.* **98**, 034101 (2007).
 - [17] R. Gutiérrez, A. Amann, S. Assenza, J. Gómez-Gardenes, V. Latora, and S. Boccaletti, *Phys. Rev. Lett.* **107**, 234103 (2011).
 - [18] L. Prignano and A. Diaz-Guilera, *Phys. Rev. E* **85**, 036112 (2012).
 - [19] S. Fortunato, *Phys. Rep.* **486**, 75 (2010).
 - [20] J. Tegnér, M. K. S. Yeung, J. Hasty, and J. J. Collins, *Proc. Natl. Acad. Sci. U.S.A.* **100**, 5944 (2003).
 - [21] F. Varela, J.-P. Lachaux, E. Rodriguez, and J. Martinerie, *Nat. Rev. Neurosci.* **2**, 229 (2001).
 - [22] R. Smith, S. Guyomarch, T. Mandel, D. Reinhardt, C. Kuhlemeier, and P. Prusinkiewicz, *Proc. Natl. Acad. Sci. U.S.A.* **103**, 5 (2006).
 - [23] I. Ruvinsky and J. J. Gibson-Brown, *Development* **127**, 5233 (2000).
 - [24] J. E. Rosenthal and G. M. Murphy, *Rev. Mod. Phys.* **8**, 317 (1936).
 - [25] A problem is in the complexity class NP if it can be solved in polynomial time by a nondeterministic Turing machine. A problem is said to be NP-complete if it is in NP and it is NP-hard, i.e., at least as hard to solve as any other problem in NP. See Ref. [26] for an in-depth discussion of complexity classes.
 - [26] C. H. Papadimitriou, *Computational Complexity* (Addison-Wesley, Reading, MA, 1994).
 - [27] E. M. Luks, *J. Comput. Syst. Sci.* **25**, 42 (1982).
 - [28] P. Holme, *Phys. Rev. E* **74**, 036107 (2006).
 - [29] P. Holme, *J. Korean Phys. Soc.* **50**, 300 (2007).
 - [30] A. Bergner, M. Frasca, G. Sciuto, A. Buscarino, E. J. Ngamga, L. Fortuna, and J. Kurths, *Phys. Rev. E* **85**, 026208 (2012).
 - [31] H. Aufderheide, L. Rudolf, and T. Gross, *New J. Phys.* **14**, 105014 (2012).
 - [32] G. Russo and J.-J. E. Slotine, *Phys. Rev. E* **84**, 041929 (2011).
 - [33] I. Fischer, R. Vicente, J. M. Buldú, M. Peil, C. R. Mirasso, M. C. Torrent, and J. García-Ojalvo, *Phys. Rev. Lett.* **97**, 123902 (2006).
 - [34] R. Vicente, L. L. Gollo, C. R. Mirasso, I. Fischer, and G. Pipa, *Proc. Natl. Acad. Sci. U.S.A.* **105**, 17157 (2008).
 - [35] A. Viriyopase, I. Bojak, M. Zeitler, and S. Gielen, *Front. Comput. Neurosci.* **6**, 49 (2012).
 - [36] In general, Eq. (1) do not admit a fully synchronized stationary state, because the system $\Omega = \dot{\theta}_i = \omega - \lambda k_i \sin \alpha \forall i$ has no solutions except for regular graphs, where $k_i = k \forall i$.
 - [37] Each row and each column of a permutation matrix has exactly one entry equal to one and all others equal to 0. If P is a permutation matrix, the matrix product PA swaps pairs of rows of A , while AP^{-1} swaps pairs of columns of A .
 - [38] See Supplemental Material at <http://link.aps.org/supplemental/10.1103/PhysRevLett.110.174102> for (i) the derivation of the expected order parameter for sets of two and seven incoherent oscillators; (ii) details on the procedure employed to compute the maximum Lyapunov exponent of Eq. (1); (iii) further information about brain data acquisition and processing.
 - [39] D. M. Abrams and S. H. Strogatz, *Phys. Rev. Lett.* **93**, 174102 (2004).
 - [40] C. R. Laing, *Physica (Amsterdam)* **238**, 1569 (2009).
 - [41] M. Shanahan, *Chaos* **20**, 013108 (2010).
 - [42] E. A. Martens, C. R. Laing, and S. H. Strogatz, *Phys. Rev. Lett.* **104**, 044101 (2010).
 - [43] Y. Iturria-Medina *et al.*, *Cereb. Cortex* **21**, 56 (2011).